Comparison of analog and digital Fourier transforms in medical image analysis

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1 Introduction

Digital processing of medical images and computer-aided diagnosis (CAD) have been studied extensively in order to improve the sensitivity and specificity of radiography for clinical diagnosis.1–3 However, radiographic images contain large amounts of information. A state-of-the-art film digitizer may create digital images with dimensions of $4K \times 5K \times 12$ bit (40 megabytes); a prototype system for full-field digital mammography acquires digital mammograms in the $4K \times 6K \times 12$ bit format.4–7 Digital computers process image data in series, therefore requiring either fast computers or long processing times. Without future development that would dramatically increase the speed and reduce the cost, use of the current digital processing techniques can only be applied to certain select cases in academic institutions; it may only be suitable for images with a small number of pixels or for procedures that do not require real-time processing. On the other hand, optical processing manipulates image data in parallel, thus the processing speed is not a limiting factor. Of course, lens aberrations and vignetting affect the accuracy of optical processing. Therefore, it is important to understand the advantages and disadvantages of both techniques for two-dimensional image processing in order to select the right method for medical image analysis and clinical diagnosis.

Abstract. The effectiveness and limitations of medical image processing using analog and digital methods are studied. Several types of errors introduced during the image processing are analyzed. For the analog optical Fourier transform, errors are introduced by the vignetting effect and lens aberration. For the digital Fourier transform, errors are introduced by the aliasing effect and the band limit. To compare the results obtained by the two techniques, a set of x-ray images was processed both optically and digitally. The former was achieved by an optical system containing a large Fourier telephoto lens and the latter by a personal computer using a Fourier transform algorithm. The veracity of both the optical and digital Fourier spectra is analyzed. Our results indicate that the optical method has high speed due to parallel processing. High veracity can be achieved in high frequency regions by using an optimal optical system. In comparison, the digital method has the advantages of high processing precision and programmability, but has low processing speed. The comparison of the two different techniques presented in this article can provide a basis for selection of the processing method in different clinical settings. Even with today’s fast computers, the optical method is still suitable for many clinical applications. The best choice lies in an analog–digital combination.

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Two-dimensional (2D) image processing has been shown to be a powerful technology in many fields, such as in pattern recognition, climate reconnaissance, ocean monitoring, and medical diagnoses.8–10 In general, there are two ways to process a 2D image. The first is to process the 2D images directly in the spatial domain. Image processing in the spatial domain can be performed using a variety of methods, such as the mathematical morphological image processing, statistical image processing, differential image processing, and fuzzy logic image processing.11–13 The second is to process the images in the frequency domain. Mathematical operations are used to transform an image into a frequency distribution and processing is performed in the frequency domain. Next, the processed data is transformed into the spatial domain again through the inverse mathematical transform. The mathematical transformation and frequency domain processing can also be performed using different methods, such as Fourier transform, Hadamard transform, Walsh transform, hoteling transform, and wavelet transform image processing.14

Image processing in the spatial domain is straightforward and simpler than in the frequency domain, because no mathematical transformations between the spatial domain and frequency domain are needed.15 Spatial-domain processing usually involves operations such as scale and intensity refinements of the input images. Certain applications such as pseudocolor density encoding and phase-only filtering for 2D...
images cannot be processed in the spatial domain; such applications require processing in the frequency domain. Image processing in the frequency domain corresponds to manipulation of the frequency information of images.

Among the mathematical transformations used in 2D image processing, the Fourier transform is one of the most important techniques. Optical Fourier transform can be divided into two diffraction processes during image processing, according to Abbe theory. First, the light falls normally upon a diffraction grating (the object), and several orders of diffraction are produced and converged to a series of points in the frequency domain. Second, since the diffraction orders are continuous, the intensities of their inverse Fourier transforms will overlap in the spatial domain. Therefore, the relationship between the object and image can be formulated as a double-Fourier transform. The exact image can be obtained only if the Fourier transform lenses are aberration free and if all the diffraction orders of the input object can be recovered during Fourier transform. The veracity of the 2D Fourier transform will influence the output results, especially in the case of medical imaging.

The main advantage of the optical (analog) Fourier transform is its speed. The input image is represented by a 2D intensity function and the intensities of all pixels of the input images are processed in parallel. Therefore, the analog method is suitable for real-time processing. The veracity of Fourier transform mainly depends on the aberration and the size of the Fourier transform lens. If the aberrations are corrected and the aperture-limitation effect of the lens is minimized, the veracity of 2D Fourier transform could be ideal. However, it is impossible to achieve such perfection because of the existence of aberrations and limitations of the aperture size of all lenses.

Different from the analog method, Fourier transform using a computer needs to digitize the input image first. The input image can also be represented by a 2D density function, but in this case each complex number will be represented by a binary series at a certain point, and then complete set of points distribution can be used as representation of the image. At the same time, some warping errors exist during image distribution processing. Furthermore, the aliasing effect, which is produced during sampling processing, will strongly influence the veracity of digital Fourier transforms. However, benefits of the digital method include high accuracy, programmability, and flexibility. Furthermore, digital processes can be easily modified in order to implement different kinds of functions and algorithms. Both methods, optical (analog) and digital (computer) Fourier transforms, can be used for 2D image processing. Each method has its own advantages and disadvantages.

In this article, the errors in both optical and digital Fourier transforms are discussed. To compare the two methods experimentally, a set of x-ray images was processed using both optical and digital methods and the results are analyzed. The selection of the processing method in real applications, based on different limiting factors, is also discussed in this article.

2 Properties and Errors of the Optical Fourier Transform

2.1 Optical Fourier Transform Properties of Lenses

The Fourier transform of a two-dimensional image can be implemented using a positive lens in an optical system. Such a system is shown in Figure 1(a), where the object plane $P_1$ serves as a light source, $P_2$ is the back focal plane of the Fourier transform lens, and $P_3$ is the principal plane of the thin Fourier transform lens. For a given complex light field $f(x,y)$ on $P_1$, its conjugate complex light field $G_f(\alpha,\beta)$ on $P_2$ can be determined by the Huygens’s principle

\[
G_f(\alpha,\beta) = C \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \left[ f(x,y) \exp(ikr') \right] dx dy \times T(p,q) \exp(ikr'') dp dq,
\]

where $C$ is a constant, and $T(p,q)$ is the transmittance of the lens.
where \( S_1 \) and \( S_3 \) denote the \( P_1 \) and \( P_3 \) surface integrals, respectively, \( C \) is an arbitrary complex constant, and \( k \) \((=2\pi\lambda)\) is the wave number. The distance \( r' \) between the point \((x,y)\) on \( P_1 \) and the point \((p,q)\) on \( P_2 \) and the distance \( r'' \) between \((p,q)\) and \((\alpha,\beta)\) on \( P_2 \) are defined as

\[
\begin{align*}
    r' &= [L^2 + (p-x)^2 + (q-y)^2]^{1/2} \\
    &= L + \frac{1}{2L}([p-x]^2 + [q-y]^2) \\
    r'' &= [f^2 + (\alpha-p)^2 + (\beta-q)^2]^{1/2} \\
    &= f + \frac{1}{2f}[(\alpha-p)^2 + (\beta-q)^2],
\end{align*}
\]

where \( L \) is the distance between \( P_1 \) and \( P_3 \), and the distance between \( P_2 \) and \( P_3 \) is the focal length \( f \) of the Fourier transform lens. The phase distribution function of the thin Fourier transform lens is defined as \( T(p,q) = \exp[-i(\pi\lambda)/(p^2 + q^2)] \).\(^{21}\) Note that the approximation in Eq. (2) is taken under the assumption that the perpendicular distance \( L \) between \( P_1 \) and \( P_3 \) and the focal length \( f \) of the lens are both much larger than the dimensions of the input and output images. Ignoring the spatial quadratic phase variation, one can obtain\(^{22}\)

\[
G_L(\alpha,\beta) = C_1 \exp\left[-i\frac{\pi}{\lambda f}\left(\frac{1-v}{v}\right)(\alpha^2 + \beta^2)\right] \times \int_{S_1} f(x,y) \exp\left[-i\frac{k}{f}(\alpha x + \beta y)\right] dxdy,
\]

where \( v = f/L \) and \( C_1 \) is an arbitrary complex constant.

Equation (3) shows that \( G_L(\alpha,\beta) \) is the 2D Fourier transform of \( f(x,y) \). In fact, if \( L = f \) the quadratic phase factor in Eq. (3) equals to unity. Thus Eq. (3) can be written as

\[
G_L(\alpha,\beta) = C_1 \int_{S_1} f(x,y) \exp\left[-i\frac{k}{f}(\alpha x + \beta y)\right] dxdy.
\]

It must be emphasized that the exact Fourier transform takes place under the condition \( L = f \) only. In other words, the input image must be placed in the front focal plane of the lens. However, errors are inevitable due to the well-known vignetting effect\(^{23}\) and to aberrations of the Fourier lens.\(^{24}\) Such errors are analyzed next.

### 2.2 Vignetting Effect of the Lens and Its Influence on the Optical Fourier Transform

Due to the limitation in the size of the lens aperture, the high-frequency components of the Fourier transform spectrum vanish using such an optical system.\(^{23}\) In order to illustrate and to analyze the impact of the vignetting effect, two basic concepts need to be introduced: the point-spread function (PSF) of the lens and the coherent transfer function (CTF) of the optical system.\(^{25}\)

In an optical system, each point on the object plane can be considered a point light source. The power of the point source can be regarded as the superposition of the powers of surrounding pixels in a small area. Furthermore, the optical intensity at each point on the output image plane is contributed to by all the points on the input object plane when the image passes through the lens. Therefore, the optical intensity on the image plane caused by a given point on the object plane point has a spatial distribution, as illustrated in Figure 1(b). Such an intensity distribution is defined as the PSF of the lens, denoted as \( h(x,y;\alpha,\beta) \). Thus, the total optical intensity distribution at any given point on the image plane can be expressed as

\[
g(\alpha,\beta) = \int_{S_1} f(x,y) h(x,y;\alpha,\beta) dxdy.
\]

Considering the space invariance of the amplitude mapping in the coherent optical system, Eq. (5) can be rewritten as\(^{25}\)

\[
g(\alpha,\beta) = \int_{S_1} f(x,y) h(\alpha-x,\beta-y) dxdy.
\]

Using the convolution theorem, one can obtain the Fourier transform of \( g(\alpha,\beta) \) as

\[
G(u,v) = \mathcal{F}\{g(\alpha,\beta)\}
\]

\[
= \int_{S_2} \exp\left[-i(u\alpha + v\beta)\right] g(\alpha,\beta) d\alpha d\beta
\]

\[
= F(u,v) \times H(u,v),
\]

where \( F(u,v) \) is the Fourier transform of \( f(x,y) \) and \( H(u,v) \) is the Fourier transform of \( h(x,y) \):

\[
\begin{align*}
H(u,v) &= \mathcal{F}\{h(x,y)\}
\end{align*}
\]

\[
= \int_{S_1} h(x,y) \exp\left[-i(ux + vy)\right] dxdy
\]

\[
= \frac{G(u,v)}{F(u,v)}.
\]

The function \( H(u,v) \), which equals the ratio of the Fourier spectrum of optical intensity on the image plane to that of the object plane, is called CTF. It represents the veracity of the Fourier transform of the coherent optical system.

Calculation of the CTF under the condition of limited aperture of the lens illustrates the frequency response of the coherent imaging system. As shown in Figure 1(b), the aperture of the lens is the exit pupil of the total optical system; the distance between the exit pupil and the back focal plane is \( d_1 \). For a converging lens of diameter \( D \), CTF, \( H(u,v) \), can be expressed as a circ function,\(^{26}\)

\[
H(u,v) = \text{circ}\left(\frac{\sqrt{u^2 + v^2}}{f_0}\right) = \begin{cases} 
1 & \sqrt{u^2 + v^2} \leq f_0 \\
0 & \sqrt{u^2 + v^2} > f_0 
\end{cases}
\]

where \( f_0 = D/2\lambda d_1 \) is a cutoff frequency of the system.

Equation (9) illustrates that the optical system with an aperture \( D \) is band limited in the frequency domain. Only the optical signal below the maximum spatial frequency \( f_s \) is allowed to pass through the lens aperture. The intensity measured within the aperture represents the squared modulus of the image’s Fourier spectrum, and a spectrum of spatial fre-
quency higher than the cutoff frequency will vanish, in spite of the fact that the image has nonzero Fourier components of high frequency.

2.3 Effect of Lens Aberration on the Fourier Transform

Because it is impossible to design and to manufacture aberration-free lenses, the error in optical Fourier transform caused by the aberration is inevitable. The effect of aberrations can be considered as the windage between the wave front out of the exit pupil and an ideal spherical wave within the aperture. Furthermore, we can assume that a phasing-shifting plate within the aperture is produced due to the aberration. The input image can be divided into many pixels, and the aberration of each pixel for the space-invariant system can be calculated using the same method as that discussed in Sec. 2.2 with a correction factor corresponding to the phase shift effect of all the pixels. Thus, the generalized coherent transfer function of the optical system can be expressed as

$$H(u, v) = \text{circ} \left( \frac{\sqrt{u^2 + v^2}}{f_0} \right) \exp[ikW(u, v)]$$

$$= \left\{ \begin{array}{ll}
\exp[ikW(u, v)] & \sqrt{u^2 + v^2} \leq f_0 \\
0 & \sqrt{u^2 + v^2} > f_0
\end{array} \right.,$$

(10)

where $W(u, v)$ is an effective path-length error due to aberration of the system, and $f_0 = D/2d_i$ is the same cutoff frequency as that in Eq. (9) for the optical system due to limitation in the size of the lens. Therefore, the band limit of the coherent optical system is not affected by the presence of the aberration.

3 Properties and Errors of the Digital Fourier Transform

When using a computer for digital Fourier transforms, the integration in Eq. (4) must be rewritten as a dispersed function.27 For a continuous input function $f(x, y)$, $f_s(x, y)$ can represent its sampling function:

$$f_s(x, y) = f(x, y) \text{comb}(x/D_x) \text{comb}(y/D_y),$$

(11)

where comb() represents a comb function, and $D_x$ and $D_y$ represent the maximum sampling distance on the $x$ axis and the $y$ axis, respectively. Then the Fourier transform of the discrete function $f_s(x, y)$ becomes

$$\mathcal{F}[f_s(x, y)] = \mathcal{F}[g(x, y)] \times \mathcal{F}[\text{comb}(x/D_x) \text{comb}(y/D_y)]$$

$$= F_s(u, v).$$

(12)

We can use the comb function to represent the sampling mathematically,

$$\mathcal{F}[\text{comb}(x/D_x) \text{comb}(y/D_y)]$$

$$= \sum_{n=-\infty}^{+\infty} \sum_{m=-\infty}^{+\infty} \delta(u - n/D_x, v - m/D_y),$$

(13)

which leads to

$$F_s(u, v) = \sum_{n=-\infty}^{+\infty} \sum_{m=-\infty}^{+\infty} F(u - n/D_x, v - m/D_y).$$

(14)

We assume that $F(u, v)$ is band limited, with a bandwidth $B_x = u_{\text{max}}$ on the $x$ axis and $B_y = v_{\text{max}}$ on the $y$ axis.28 One must choose the sampling distances according to the following rule:

$$D_x \leq \frac{1}{2B_x}$$

(15)

and

$$D_y \leq \frac{1}{2B_y}. $$

(16)

The specific sampling distances $D_x$ and $D_y$ in Eq. (15), as the reciprocal of the spatial bandwidth, are called the Nyquist rate. The Nyquist rate is the largest sampling distance that can be chosen without introducing the aliasing effect.29

Following Eqs. (11)–(15), the inverse Fourier transform of the input function can be rewritten as

$$f(x, y) = \sum_{n=-\infty}^{+\infty} \sum_{m=-\infty}^{+\infty} f(n/B_x, m/B_y) \sin \left[ 2B_x \left( x - \frac{n}{2B_x} \right) \right]$$

$$\times \sin \left[ 2B_y \left( y - \frac{m}{2B_y} \right) \right],$$

(16)

The errors in the digital Fourier transform spectra are mainly caused by the discrete processing of data. If the sampling distances are larger than the reciprocal of the spatial bandwidth, the aliasing effect will appear.

4 Comparison Between Analogy and Digital Fourier Spectra

4.1 Optical Fourier Transform of Medical Images

In our experiments, the optical (analog) Fourier transform was performed using an optical system containing a Fourier transform telephoto lens. The parameters of the optical system are as follows. The distance between the exit pupil of the lens and the lens’s back focal plane $d_i$ was 458.3 mm, the lens’s focal length $f$ was 1000 mm, its relative aperture $D/f$ was 1:12, and the maximum wave aberration was 1/8 λ. The light source is a He–Ne laser with a center wavelength of 632.8 nm. The cutoff frequency $f_0$ of the system can be calculated as

$$f_0 = \frac{D}{2\lambda d_i} = 143.67/mm.$$  

(17)

For the circular aperture this cutoff frequency is uniform in all directions of the frequency plane.

4.2 Digital Fourier Transform

The digital Fourier Transform of a 2D image was performed using the computer software Interactive Data Language (IDL), a program designed for 2D image data processing.30 It consists of a scanner that acquires images from 135 photographic films. The optical intensities of the acquired image
were digitized and the Fourier transform for the image was performed numerically. The cutoff frequency for the optical system given by Eq. (17) requires a spatial frequency for the digital Fourier transform to be 120 and 80 lines/mm on the x and y axes, respectively, for the 36 mm×24 mm films. Following Eq. (15), the maximum sampling distances should be 0.004 and 0.006 mm on the x and y axes, respectively, to achieve the same band limit as that given in Eq. (17). As a consequence, the number of sampling points should be 9000 ×4000. It is necessary to require 36 Mbit memory capacity to store one image with 256 gray levels. For simplicity and for a demonstration of the aliasing effect, we selected 3000×1333 sampling points, namely, one from every nine neighboring points. This reduced the total memory for each image to 4 Mbit.

4.3 Comparison Between Optical and Digital Fourier Spectra

A set of medical images was processed using both optical and digital Fourier transform methods. The original medical images on x-ray films were optically reduced and recorded on 36 mm×24 mm photographic films, as shown in Figures 2(a), 3(a), and 4(a). Having completed the processing procedures described in Secs. 4.1 and 4.2, the optical Fourier spectra of the medical images are given in Figures 2(b), 3(b), and 4(b), and the digital Fourier spectra are given in Figures 2(c), 3(c), and 4(c).

Figure 2(a) is an image of letter E with a black background. As expected, due to the rectangular shape of the image, the spectra in both cases showed a distribution along the x axis and the y axis. The optical Fourier spectrum in Figure 2(b) shows more details of the structure than the digital spectrum in Figure 2(c) under the same display conditions.

The optical and digital Fourier spectra of an x-ray image of breast tissue complex in structure [Figure 3(a)] are given in Figures 3(b) and 3(c). Apparently, the center region (low frequency region) of both spectra has the same distribution pattern, evidenced by the close resemblance of the center bright spots. However, differences in the periphery of the spectra are seen clearly. The optical method was able to process signals of higher frequency, as shown by the wider spectrum distribution in Figure 3(b). Optical and digital Fourier spectra of an x-ray image of bones [Figure 4(a)] are shown in Figures 4(b) and 4(c). These bone images were used to show the result when the image with more sharp contrast was inputted. Both Figures 4(b) and 4(c) show similar patterns in the center region with the spectrum distribution mainly along the bisector of the first and third quadrants. Figure 4(b), however, shows more details of the structure in the high frequency region (peripheral region).

Figures 2–4 showed the following common features of the two transformation methods: (1) the results in the center regions (low-frequency region) of the spectra are nearly identical; (2) the optical method provides more details in the high frequency regions.

The shortcomings of the digital Fourier transform in this experiment are largely due to the fact that the number of sampling points in the processing is limited. If enough pixels were sampled to satisfy the same frequency cutoff like in optical method, the veracity of both transformations would be...
similar. Under the specific conditions of our experiments, the digital processing points required should be 9000 \times 4000 based on the required sampling distances given by Eq. (15). However, the digital Fourier transform in our experiments required about 30 s to process an image with 3000 \times 1333 points and it will require an estimated 500 s to process one of 9000 \times 4000 points using a personal computer. The spectra in Figures 2(c), 3(c), and 4(c) obtained using fewer sampling points—one ninth of the points required for the same optical cutoff frequency—clearly demonstrated the aliasing effect. If enough pixels were sampled, the digital method can provide highly accurate results, provided that the processing time is not a limiting factor.

The disadvantages of optical method mainly lie in the complex structure of the optical system, which requires an assembly of various optical components. Thus, the level of difficulty of data processing increases. Furthermore, all the analyses here are based on the assumption of a thin lens. In reality, the thickness of the lens will always introduce noticeable errors. Based on our analysis of errors introduced by lens aberration [see Eq. (10)], optical processing can provide the correct power of the spectrum, but the accuracy of the phase distribution is limited.

On the positive side, the optical method has the advantages of real-time parallel processing for a Fourier transform of larger size images. The computer can provide greater flexibility and high accuracy.

5 Discussions

The Fourier transform is one of the most important methods in 2D image processing. It has several properties that make it attractive for image processing applications. In this article, the characteristics of analog (optical processing) and discrete (digital processing) Fourier transforms are analyzed.

Specifically, we discussed the effect of lens aberration on the optical Fourier transform by introducing the general phase distribution function. As a consequence, the lens aberrations only contribute a phase correction factor, and do not affect the overall intensity and the cutoff frequency of the Fourier spectrum. The vignetting effect contributes to errors of the optical spectrum by introducing a cutoff frequency that is dependent on the size of the lens.

The main factors that contribute to the errors in the digital Fourier transform are the aliasing effect and the band limit. The former is due to sampling of a small number of pixels and it affects the veracity of the spectrum in the high frequency region, as shown in our experiments (see the spectra in Figures 2–4). The latter is due to the requirement of digital processing; the band limit automatically restricts the high-frequency components in the Fourier spectrum, as determined by Eq. (15). In principle, digital processing can provide a high degree of veracity if there is a large enough number of sampling points. However, a long processing time is required in such cases. Therefore, the difficult task is to keep a balance between the high accuracy desired (more pixels) and the processing speed desired (fewer pixels).

Overall, the error analysis and our experimental results indicate that the outcomes of the two Fourier transform methods for medical images are similar in the center region of the Fourier plane (the lower-frequency region), but different in the peripheral region (the higher-frequency region). In general, the analog method (optical processing) is suitable for medical images with a large number of pixels and for clinical procedures that require real-time operations, such as fluoroscopy because of its required high processing speed. Digital processing, on the other hand, can provide greater flexibility and programmability.

In summary, optical processors have good analog quality and high speed through parallel processing. In order to achieve high accuracy, optical components with minimal aberration are needed. In addition, current optical processors are less flexible since they are more limited in programmability than digital processors. One technical solution with which to improve their flexibility and programmability is to develop hybrid devices, namely, parallel optical processors equipped with optimal digital–optical interfaces.

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